

# Sterile Neutrinos I

By R.G. Van de Water  
Los Alamos National Laboratory  
July 9, 2007  
FNAL Neutrino Summer School

# Lecture Outline

- Sterile I: Theory and Phenomenology
  - Properties of sterile neutrinos in the standard model.
  - Sterile neutrinos and beyond the standard model, LRSM as an example.
  - Sterile neutrinos and extra dimensions.
  - Sterile neutrinos in cosmology/astrophysics.
  - 3+2 models, fits to the data.
- Sterile II: Measurements
  - How to make a smoking gun active-sterile neutrino oscillation measurement. A primer!

# Sterile Neutrinos, Who needs them!

- Discovery of Sterile neutrinos would rank in importance with discovery of charm, bottom, and tau.
- These discoveries was presaged by theory to solve problems of tree-level flavor changes and the need for a third generation to introduce CP violations that were observed in experiments.
- May point to physics beyond the standard model.

# Sterile Neutrinos in the Standard Model Gauge Group

(see R.R. Volkas, hep-ph/0111326)

- SM gauge group as constructed,

$$G_{SM} = SU(3)_c \otimes SU(2)_L \otimes U(1)_\gamma$$

- One generation/family of quarks and leptons forms the reducible representation,

$$Q_L = \begin{pmatrix} u_L \\ d_L \end{pmatrix} \sim (3, 2)(1/3), \quad d_R \sim (3, 1)(-2/3), \quad u_R \sim (3, 1)(4/3);$$

$$\ell_L = \begin{pmatrix} \nu_L \\ e_L \end{pmatrix} \sim (1, 2)(-1), \quad e_R \sim (1, 1)(-2), \quad ? \text{ MISSING ENTRY ?}$$

For massless neutrinos, right handed states do not exist, i.e. handedness is conserved.

$$v_L = \frac{(1-\gamma)}{2}v, \text{ helicity} = -1 \text{ for } m=0$$

$$v_R = \frac{(1+\gamma)}{2}v, \text{ helicity} = +1 \text{ for } m=0$$

- Parity violation put into SM by hand, i.e.  $(1-\gamma)$  coupling only,  $(1+\gamma)$  coupling not observed.
- For minimal SM (no neutrino mass), mismatch between quark and lepton degrees of freedom!

# Generating Dirac Neutrino Mass

- Quark and charged lepton (Dirac) mass generated through Yukawa coupling.

$$\mathcal{L}_{\text{Yuk}} = h_d \bar{Q}_L d_R \Phi + h_u \bar{Q}_L u_R \tilde{\Phi} + h_e \bar{\ell}_L e_R \Phi + \text{H.c.}$$

- With spontaneous symmetry breaking, Dirac neutrino mass terms of type,

$$m_D \bar{\nu}_L \nu_R$$

- If  $\nu_R$  are absent, and the Higgs sector remains minimal, then neutrinos are massless.
  - neutrino mass implies  $\nu_R$  exists!
- $\nu_L$  and  $\nu_R$  are building blocks of the neutrino mass Lagrangian.

# Sterile Neutrinos in the Standard Model Gauge Group

- Missing right handed neutrino state should be,  
$$\nu_R \sim (1,1)(0).$$
- where (color,weak)(hypercharge =  $2*(Q - I_L)$ ).
- $\nu_R$  has the quantum numbers of the vacuum, thus sterile with respect to the standard model gauge interactions!
- SM with neutrino mass now looks like,

$$\begin{aligned} q_L = \begin{pmatrix} u_L \\ d_L \end{pmatrix} &\sim (3,2)(1/3), & d_R &\sim (3,1)(-2/3), & u_R &\sim (3,1)(4/3); \\ \ell_L = \begin{pmatrix} \nu_L \\ e_L \end{pmatrix} &\sim (1,2)(-1), & e_R &\sim (1,1)(-2), & \nu_R &\sim (1,1)(0) \end{aligned}$$

# Sterile Neutrinos in the Standard Model Gauge Group

- Presence of  $\nu_r$  would enhance two aesthetic qualities:
  - Left-right similarity, i.e. for each left handed fermion, there is a right handed one.
  - Quark-lepton similarity, i.e. For each quark of a given chirality (handedness), there is an associated chiral lepton.
- Family structure of a quark-lepton family provides strong motivation for one sterile fermion per family.
- Historically, quark-lepton similarity used to predict charm.

# Majorana Neutrino Mass

- Because Quarks/Leptons have charge, they only have Dirac masses which are equivalent between left and right components.
- But only the neutrino can have Majorana mass terms of the type,

$$M \bar{\nu}_R^c \nu_R$$

- With the full mass matrix given by,

$$\begin{pmatrix} \bar{\nu}_L & \overline{(\nu_R)^c} \end{pmatrix} \begin{pmatrix} 0 & m_D \\ m_D & M \end{pmatrix} \begin{pmatrix} (\nu_L)^c \\ \nu_R \end{pmatrix}.$$

- $m_D$  is the Dirac mass  $\sim$  EW symmetry breaking.
- $M$  is the Majorana mass  $\sim$  unknown value.

# Sterile Neutrino Mass: See Saw Model

- If we assume  $M \gg m_D$ , eigenvalues ( $m_L$  and  $m_R$ ) become approximately,

$$\frac{m_D^2}{M} \quad \text{and} \quad M,$$

- If we assume  $M \sim$  GUT scale ( $\sim 10^{15}$  GeV), and  $m_D \sim M_{EW}$  ( $\sim 10^2$  GeV), then:

$$m_L = 10^{-2} \text{ eV}, \quad m_R = MGUT$$

- And active-sterile mixing angle  $\tan\theta \ll 1$ .
- We explain smallness of neutrino mass, and lack of active-sterile couplings.
- But why  $M \sim$  GUT scale???

# Sterile Neutrino Mass: Pseudo Dirac

- Interesting limits  $M \ll m_D$  (pseudo-Dirac), eigenvalues ( $m_L$  and  $m_R$ ) become approximately,

$$m_D \pm \frac{M}{2}$$

- As  $M \rightarrow 0$  we end up with,
  - Nearly degenerate pair with a mass gap  $m_D$  above zero, i.e.  $m_L \sim m_R \sim m_D$
- With mixing between the eigenstates given by,

$$\tan 2\theta = -\frac{2m_D}{M} \Rightarrow |\theta| \simeq \frac{\pi}{4}.$$

- Nearly maximal active(L)-sterile(R) mixing.
- Light sterile neutrinos implies large mixing!!!

# Implications of Sterile Neutrino Mass

- Thus, non observations of sterile neutrinos implies  $m_R \sim \text{MGUT}$ , which begs the question,
  - Why is the Majorana mass at such a large symmetry breaking scale??
- Observations of sterile neutrinos with small mass ( $m_R \sim M_{EW}$ ) and large coupling to active neutrinos, can be accommodated by the SM,
  - How would we explain smallness of neutrino Dirac mass?
- Searching for sterile neutrinos is important for understanding source of fermion masses and new physics!

# Sterile Neutrinos Beyond the SM

- Examples of extensions to the SM;
  - Mirror matter model (weakly sterile).
  - Pati-Salam model (weakly sterile).
  - SU(5) grand unification (fully sterile).
  - SO(10) grand unification (weakly sterile).
  - Left-right symmetric models (weakly sterile).
  - Unusual left-right symmetric (fully sterile).
- Notes,
  - “Fully sterile” feels no gauge interactions, including those beyond the SM.
  - “Weakly sterile” does not feel standard model gauge interactions (strong, EM, and weak).
  - Fully/Weakly sterile neutrinos do couple to gravity. Fully sterile can interact via Higgs exchange and partake in mass mixing.

# Sterile neutrinos beyond the SM

## Left-Right Symmetric Model

- Basic motivation of LRSM is to treat left and right handed fermions more symmetrically. Parity violation is induced spontaneously ( $> \text{few TeV}$ ) rather than engineered explicitly.
- The gauge group is now given by,

$$G_{\text{LR}} = \text{SU}(3)_c \otimes \text{SU}(2)_L \otimes \text{SU}(2)_R \otimes \text{U}(1)_{B-L},$$

- $\nu_r$  is now explicitly put into the model and participate in right handed weak interactions mediated by exotic  $W$ -like bosons and an additional  $Z'$  boson (weakly sterile).

# Sterile neutrinos beyond the SM

## Left-Right Symmetric Model

- LRSM completed by specifying the Higgs sector, which is constructed to yield a see-saw structure for neutrinos and breaks  $G_{LR}$  in two stages.

- The resulting VEV hierarchy results,

$$\langle \Delta_R \rangle \sim v_R \gg \langle \Phi \rangle \sim v_{ew} \gg \langle \Delta_L \rangle \sim v_L$$

- At the scale  $v_R$ , right handed weak isospin is spontaneously broken and the  $W_R$  and  $Z'$  acquire mass.
- This symmetry breaking automatically generates large see-saw Majorana masses within the theory (not added ad hoc) and connects with see-saw limit.
- Nice idea, but not established. Example of neutrino mass as a window onto higher symmetries.

# Neutrino Handedness

## A Question??

- Dave Griffiths “There is a danger in carrying handedness too far”, D. Griffiths, Intro to Elementary Particles (1987).
- When particles have non zero mass, handedness is not conserved in the propagation of a free particle. This means left-right handed states are not well defined.
  - Can boost to a frame where a left handed state will become right handed.
- Does this mean sterile neutrinos states are not real, or at least there mass is not different from left handed states??

# Sterile neutrinos and Extra Dimensions. A cool idea!

- See H. Pas, et al, PRD 72, 095017 (2005) for details.
- Definitions;
  - Brane -- our space-time universe
  - Bulk ---- extra dimensions. The Brane is embedded in the bulk.
- Idea, assume active-sterile neutrino mixing at the  $\sim 1$  eV<sup>2</sup> region (LSND oscillations).
- Sterile neutrinos have the same gauge quantum numbers of the vacuum and the graviton.
- Theories which include gravity usually involve extra dimensions. Thus, sterile neutrinos can in principle follow the graviton into the extra dimensions.

# Sterile neutrinos and Extra Dimensions: Brane Fluctuations

- Theories show that branes embedded in higher dimensional spacetime possesses fluctuations or buckles on a microscopic scale from gravity, thermal and/or quantum fluctuations.
- Construct simple 1 + 1 dimensional space time embedded in a 1+2 space-time “toy” model with Minkowski metric,

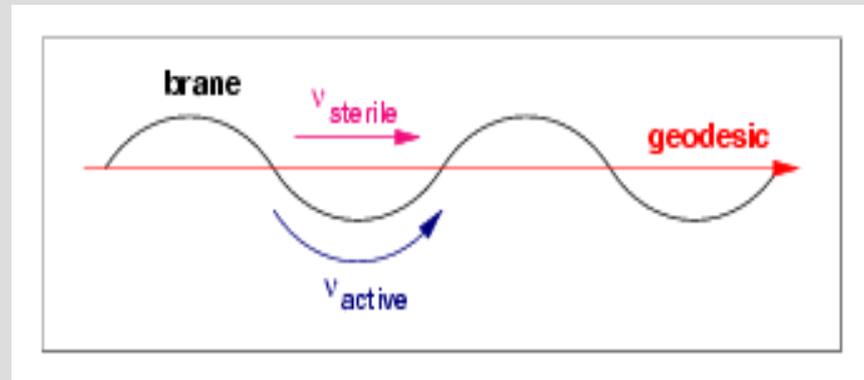
$$ds^2 = dt^2 - dx^2 - dy^2.$$

- Assume the brane exhibits the following spatial variations (no claim to realism),

$$y = A \sin kx;$$

# Sterile neutrinos and Extra Dimensions: Geodesics

- Schematically...



- With travel distance for the sterile neutrino of,

$$D_g = x.$$

- And the geodesic for active states on the brane,

$$D_b = \int_{\text{brane}} \sqrt{dx^2 + dy^2} = \int^x \sqrt{1 + A^2 k^2 \cos^2 kx} dx.$$

# Sterile neutrinos and Extra Dimensions: Shortcuts

- After some math, the shortcut distance is given by,

$$\epsilon = \left(\frac{Ak}{2}\right)^2.$$

- Where A is the amplitude and k the wave number of brane fluctuation, with Ak being the aspect ratio.
- The neutrino evolution equation in the brane is given by,

$$i \frac{d}{dt} \begin{pmatrix} \nu_a(t) \\ \nu_s(t) \end{pmatrix} = H_F \begin{pmatrix} \nu_a(t) \\ \nu_s(t) \end{pmatrix},$$

- The shorter distance traveled by the sterile neutrino will be added in as an effective potential,

$$H_F = + \frac{\delta m^2}{4E} \begin{pmatrix} -\cos 2\theta & \sin 2\theta \\ \sin 2\theta & \cos 2\theta \end{pmatrix} + E \frac{\epsilon}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}.$$

# Sterile neutrinos and Extra Dimensions: Energy Resonance

- Similar to MSW oscillation except more pronounced energy dependence with  $\Delta m^2$  varying as  $E^2$ .
- Bulk term beats against the brane term to give resonant mixing at energy,

$$E_{\text{res}} = \sqrt{\frac{\delta m^2 \cos 2\theta}{2\epsilon}}$$

- Assuming  $\epsilon \ll 1$ , then  $\Delta m^2 \ll E_{\text{res}}^2$
- Oscillation probability given by,

$$P_{\alpha\beta} = \sin^2 2\tilde{\theta} \sin^2(\delta HD/2),$$

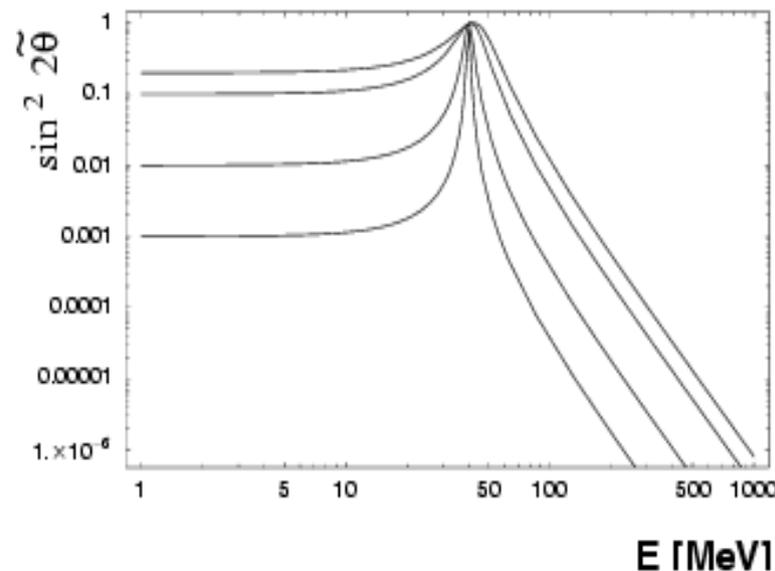
- With the new values given in terms of standard ones,

# Sterile neutrinos and Extra Dimensions: Energy Resonance

$$\sin^2 2\tilde{\theta} = \frac{\sin^2 2\theta}{\sin^2 2\theta + \cos^2 2\theta \left[ 1 - \left( \frac{E}{E_{\text{res}}} \right)^2 \right]^2},$$

$$\delta H = \frac{\delta m^2}{2E} \sqrt{\sin^2 2\theta + \cos^2 2\theta \left[ 1 - \left( \frac{E}{E_{\text{res}}} \right)^2 \right]^2}.$$

- For  $E_{\nu} < E_{\text{res}}$ , standard oscillations
- For  $E_{\nu} = E_{\text{res}}$ , resonant oscillations
- For  $E_{\nu} > E_{\text{res}}$ , oscillations die away.



$$E_{\text{res}} = 40 \text{ MeV}$$

# Sterile neutrinos and Extra Dimensions. Pre-MB predictions!

- Short baseline experiments, values of  $E_{\text{res}}$  allowed.

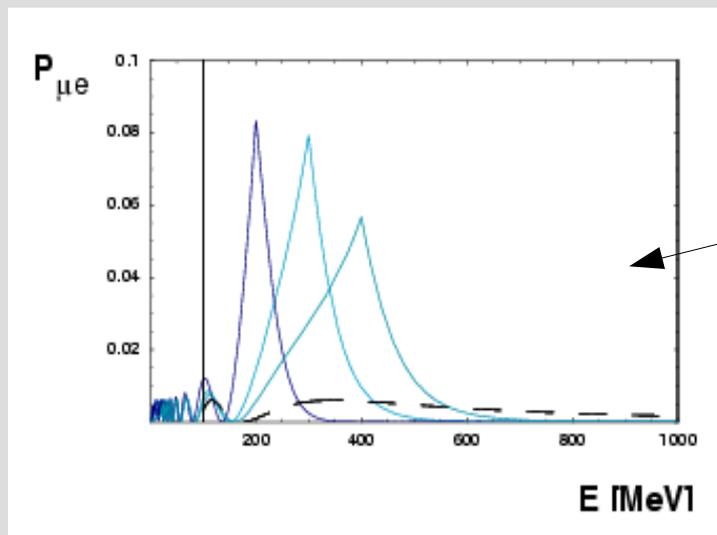
	$\alpha\beta$	$E_\nu$	$D$	$\sin^2 2\tilde{\theta}_{\alpha\beta}$
LSND	$\mu e$	20–52.8 MeV	30 m	$>0.003$
KARMEN	$\mu e$	20–52.8 MeV	17.7 m	$<0.002$
MiniBooNE	$\mu e$	0.1–1.0 GeV	540 m	$\sim 0.0006$
BUGEY	$e\bar{\nu}$	1–6 MeV	25 m	$<0.15$
CDHS	$\mu\bar{\nu}$	$>1$ GeV	755 m	$<0.1$

Standard oscillations sets lower limits above  $\sim 50$  MeV.

oscillation suppressed below  $\sim 400$  MeV

$$\sin^2 2\tilde{\theta}_{\mu\mu} \simeq \cos^2 \theta_*, \quad \sin^2 2\tilde{\theta} \simeq \cos^2 \theta_* \tan^2 2\theta \left( \frac{E}{E_{\text{res}}} \right)^{-4}.$$

- Predictions for MinibooNE,



No high energy oscillations, enhanced oscillations at low energy (100-400MeV)!

Needless to say, extraD authors are working on fits to MiniBooNE low energy excess.

# Sterile neutrinos in Cosmology

- Light sterile neutrinos can have important effects in big bang nucleosynthesis (BBN), the process thought responsible for generating the light isotopes  $^4\text{He}$ ,  $^3\text{He}$ , D, and  $^7\text{Li}$ .
- BBN epoch occurs shortly after neutrinos decouple from e/gamma plasma at  $T \sim 1 \text{ MeV}$ .
- The plasma contain nucleonic contamination which are being intra-converted through the process,

$$\nu_e n \leftrightarrow e^- p, \quad \bar{\nu}_e p \leftrightarrow e^+ n, \quad n \leftrightarrow p e^- \bar{\nu}_e.$$

$$\text{with } n/p = \exp\left[\frac{(m_p - m_n)}{T}\right] \text{ for zero lepton number}$$

- Ratio of elements determined by relative rates of above reactions, and expansion rate of the universe during relevant period.

# Sterile neutrinos in Cosmology

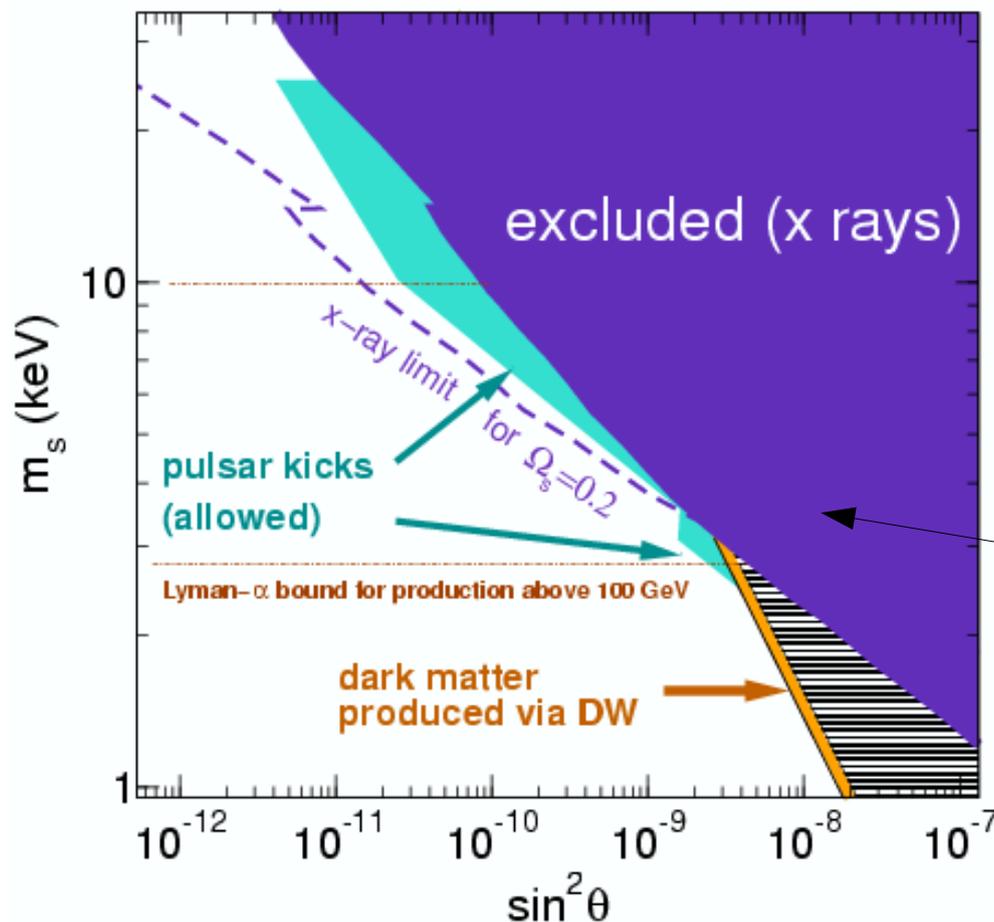
- Expansion during BBN determine by the relativistic component of the plasma. In standard BBN, the relativistic species are active (anti)neutrinos, electrons, positrons, and photons.
- A significant light sterile neutrino component would increase the expansion rate. Thus, decreasing the weak freeze-out period, increasing  $T$ , which changes the  $n/p$  ratio, resulting in an unacceptably high  ${}^4\text{He}$  yield.
- Have assumed neutrino-antineutrino number density asymmetry is zero (zero lepton number).
  - However, active-sterile oscillations can induce large neutrino flavor asymmetries.

# Sterile neutrinos in Cosmology

- Current BBN without sterile neutrinos work. However, the presence of sterile neutrinos can act to alter the  $n/p$  ratio.
  - There are many factor that can act to increase or decrease the effects.
- Cosmology with light sterile neutrinos requires careful analysis, with the outcome determined by the model and choice of parameters.
  - **Successful cosmologies can result!**

# Sterile Neutrinos in Astrophysics

- Dark matter candidate.
- Pulsar Kicks ( $\nu_s$  carries away energy asymmetry)



Limits from various astrophysics sources:

- pulsar kicks.
- relic sterile neutrino decays produce x-rays.
- $\nu_s$  dark matter below  $\Omega_s$  dotted line.

$m_s \sim$  few keV, and mixing  $\sim 10^{-10}$

# Sterile Neutrinos in Astrophysics

- Super Nova explosions (R-process).

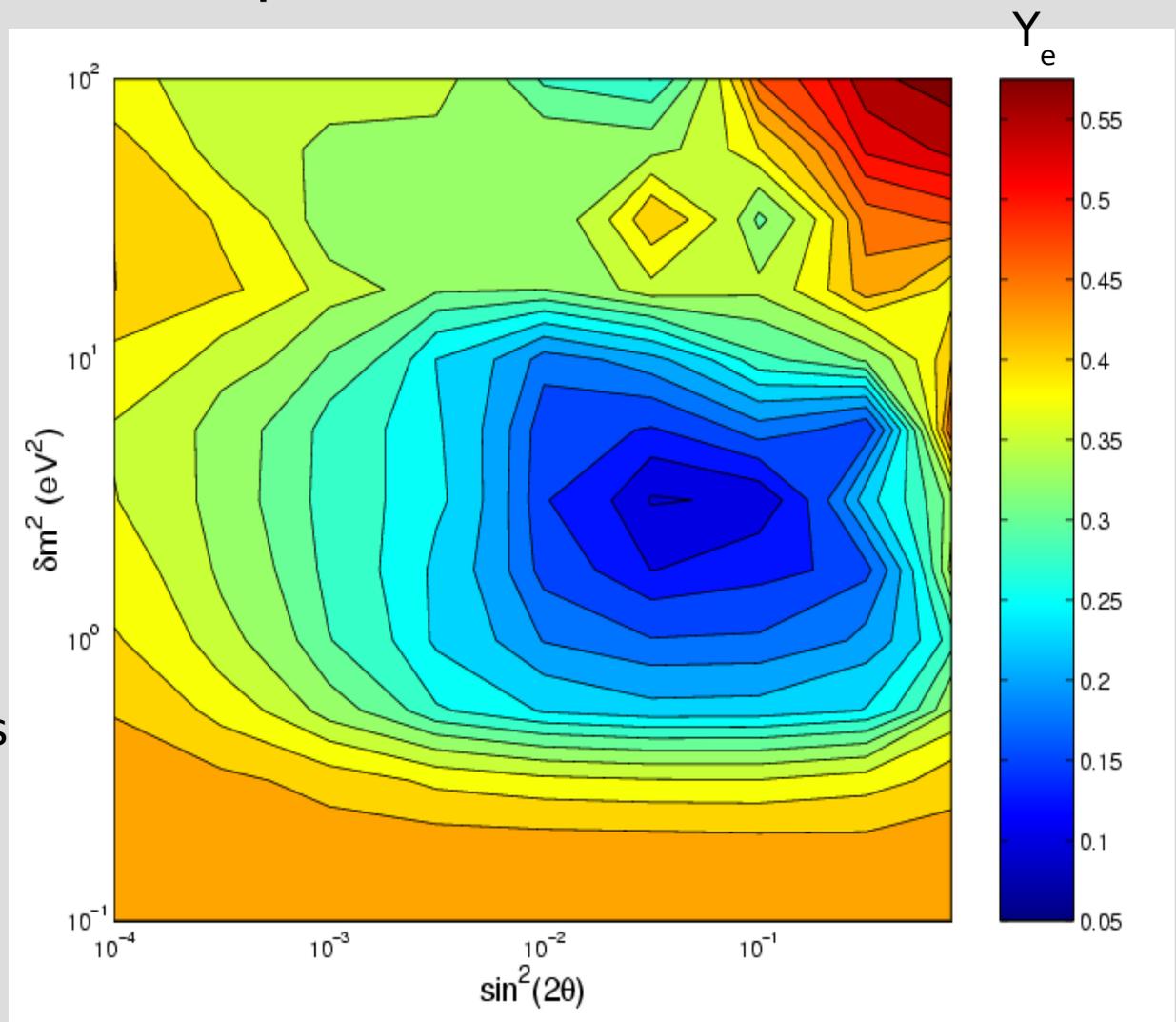
Reactions in the neutrino wind,



with  $Y_e = 1/(1+n/p)$

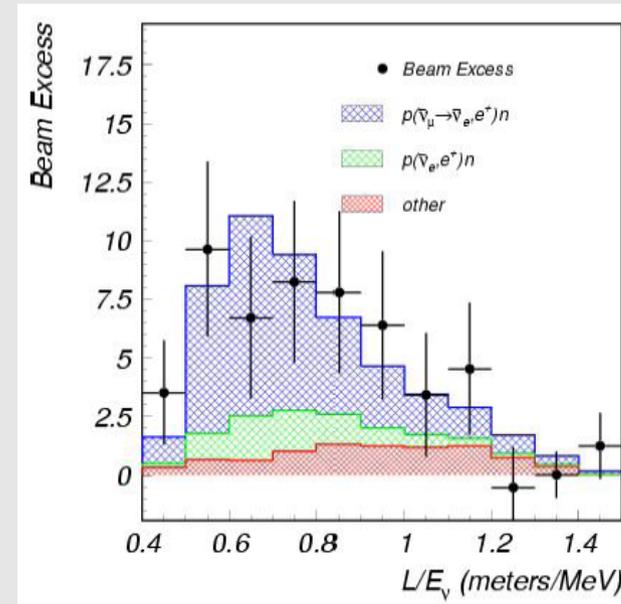
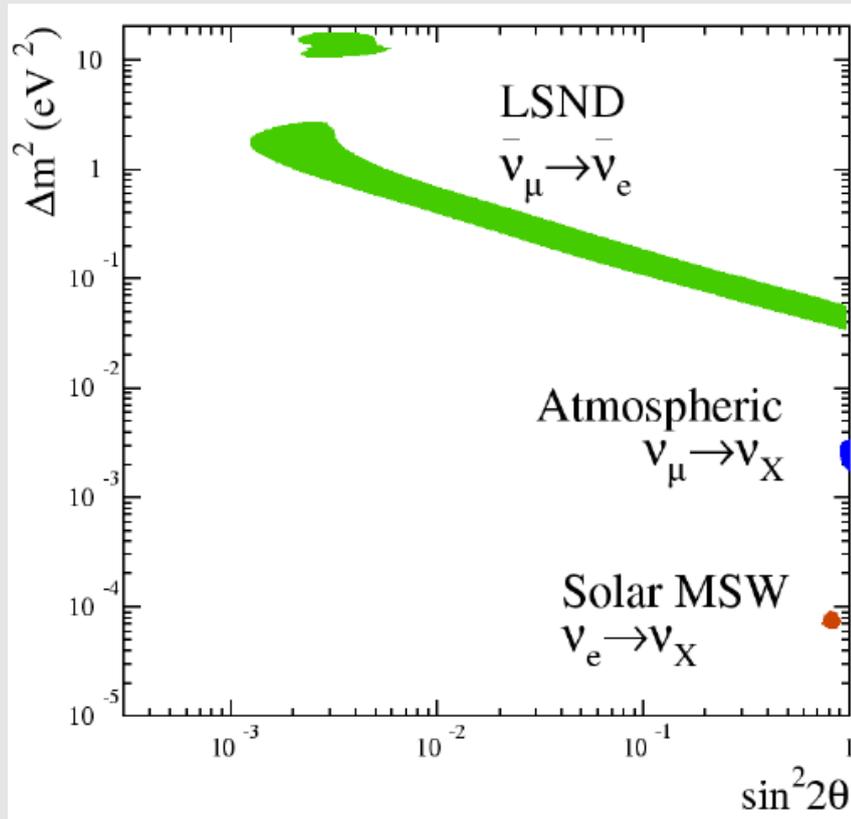
- Values of  $Y_e < 0.3$  yield neutron rich environment which is required for heavy element production.
- Presence of sterile neutrinos can reduce  $Y_e$ .

Neutrino background not included.



# Terrestrial Hints for Sterile Neutrinos

## The LSND Signal



$\Delta m^2_{\text{LSND}} \sim 0.1\text{-}10 \text{ eV}^2 + \text{small mixing}$

$\Delta m^2_{\text{LSND}} \gg \Delta m^2_{\text{atm}} \gg \Delta m^2_{\text{sol}}$

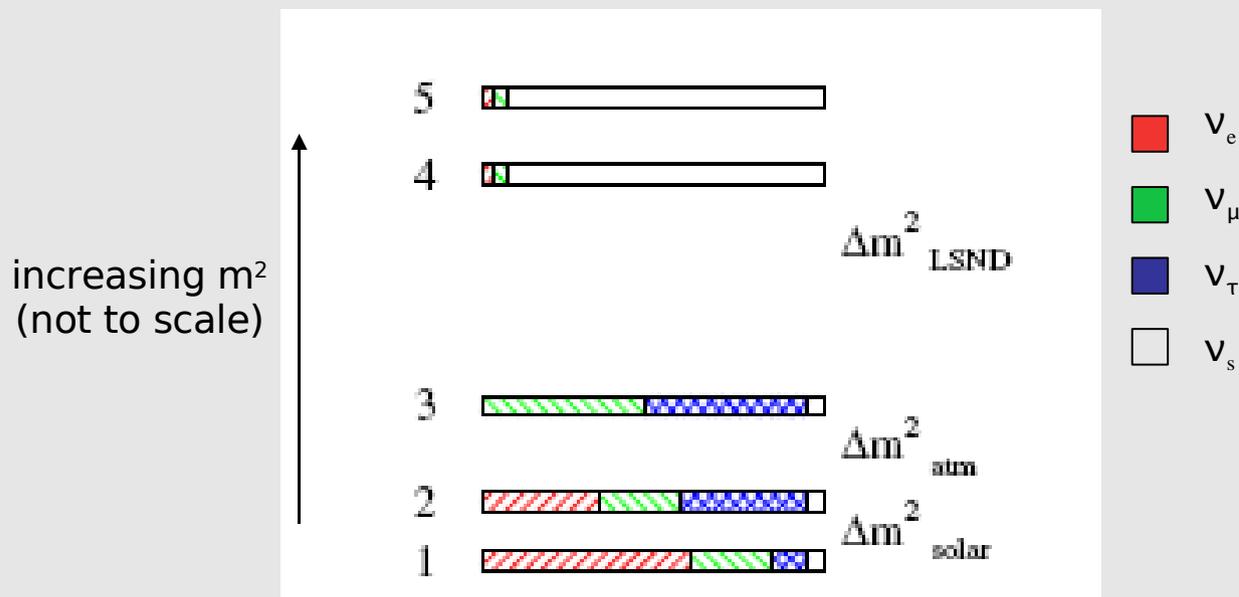
→one option: **3 active + n “sterile” neutrinos**

→other options: neutrino decay, extra-D, etc.

# 3+2 Model Phenomenology

## 3 active + 2 **sterile** neutrinos

- **light** sterile neutrinos
- they can interact thru non-standard weak couplings
- they have very small active flavor content ( $U_{e4}, \dots, U_{e5}, \dots$ )  
 → can participate in neutrino oscillations



Why  $n=2$ ?

3+1 models: SBL and LSND marginally consistent with each other

3+2: next natural step...

[M. Sorel, et al. hep-ph/0305255]

# The Mass Matrix: Adding One Sterile Neutrino

(see R.R. Volkas, hep-ph/0111326)

- In absence of Higgs triplets, the general 3+1 neutrino mass matrix is given by,

$$\begin{pmatrix} 0 & 0 & 0 & m_1 \\ 0 & 0 & 0 & m_2 \\ 0 & 0 & 0 & m_3 \\ m_1 & m_2 & m_3 & M \end{pmatrix}$$

- But this has two zero and two non-zero eigenvalues
  - Only two distinct  $\Delta m^2$  values.
  - Cannot fundamentally motivate 3+1 models to explain LSND anomaly.

# 3+2 Analysis

Idea: If light sterile neutrinos ( $\nu_s$ ) exist, then:

Includes CP phase;  $\phi = -\phi$  for antineutrinos

$\nu_\mu \rightarrow \nu_s \rightarrow \nu_e$

$$P(\nu_\alpha \rightarrow \nu_\beta) = 4|U_{\alpha 4}|^2|U_{\beta 4}|^2 \sin^2 x_{41} + 4|U_{\alpha 5}|^2|U_{\beta 5}|^2 \sin^2 x_{51} + 8|U_{\alpha 5}||U_{\beta 5}||U_{\alpha 4}||U_{\beta 4}| \sin x_{41} \sin x_{51} \cos(x_{54} - \phi_{54})$$

$\nu_\mu \rightarrow \nu_s$

$\nu_e \rightarrow \nu_s$

$$P(\nu_\alpha \rightarrow \nu_\alpha) = 1 - 4[(1 - |U_{\alpha 4}|^2 - |U_{\alpha 5}|^2)(|U_{\alpha 4}|^2 \sin^2 x_{41} + |U_{\alpha 5}|^2 \sin^2 x_{51}) + |U_{\alpha 4}|^2|U_{\alpha 5}|^2 \sin^2 x_{54}]$$

With SBL approximation  $\Delta m_{\text{solar}} = 0$ ,  $\Delta m_{\text{ATM}} = 0$ , and  $x_{ij} = \delta m_{ij} L / 4E$

**Experimental constraints from:**  
**LSND, KARMEN, NOMAD, MB, CCFR, CDHS, CHOOZ, BUGEY (+ atm constraint)**

appearance experiments  
 $(\nu_\mu \rightarrow \nu_e)$

disappearance experiments  
 $(\nu_\mu \rightarrow \bar{\nu}_\mu \text{ or } \nu_e \rightarrow \bar{\nu}_e)$

$(\nu_\mu$  disappearance  
 Constraint)

3+2 models can produce differences between neutrino and antineutrino appearance rates!

# SBL Combined Analysis

Short-baseline (SBL) experiments on  $\nu_\mu$  and  $\nu_e$  disappearance, and on  $\nu_\mu \rightarrow \nu_e$  appearance, probe the same  $\Delta m^2$  range and matrix elements

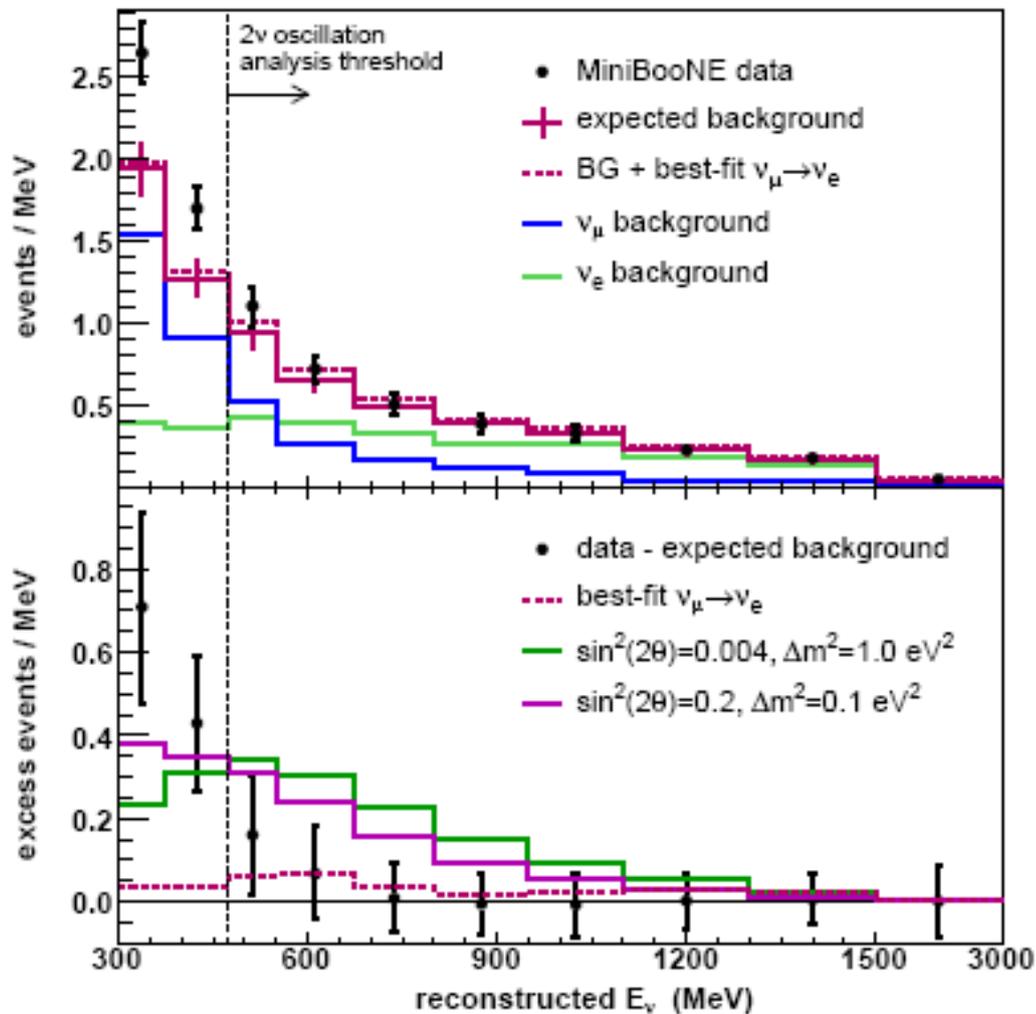
$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \\ \nu_s \\ \nu_{s'} \\ \vdots \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} & U_{e4} & U_{e5} & \dots \\ U_{\mu1} & U_{\mu2} & U_{\mu3} & U_{\mu4} & U_{\mu5} & \dots \\ U_{\tau1} & U_{\tau2} & U_{\tau3} & U_{\tau4} & U_{\tau5} & \dots \\ U_{s1} & U_{s2} & U_{s3} & U_{s4} & U_{s5} & \dots \\ U_{s'1} & U_{s'2} & U_{s'3} & U_{s'4} & U_{s'5} & \dots \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \\ \nu_4 \\ \nu_5 \\ \vdots \end{pmatrix}$$

Channel	Experiment	Lowest $\Delta m^2$ Reach (90% CL)	$\sin^2 2\theta$ Constraint (90% CL)	
			High $\Delta m^2$	Optimal $\Delta m^2$
$\nu_\mu \rightarrow \nu_e$	LSND	$3 \cdot 10^{-2}$	$> 2.5 \cdot 10^{-3}$	$> 1.2 \cdot 10^{-3}$
	KARMEN	$6 \cdot 10^{-2}$	$< 1.7 \cdot 10^{-3}$	$< 1.0 \cdot 10^{-3}$
	NOMAD	$4 \cdot 10^{-1}$	$< 1.4 \cdot 10^{-3}$	$< 1.0 \cdot 10^{-3}$
$\nu_e \rightarrow \nu_{s'}$	Bugey	$1 \cdot 10^{-2}$	$< 1.4 \cdot 10^{-1}$	$< 1.3 \cdot 10^{-2}$
	CHOOZ	$7 \cdot 10^{-4}$	$< 1.0 \cdot 10^{-1}$	$< 5 \cdot 10^{-2}$
$\nu_\mu \rightarrow \nu_\mu$	CCFR84	$6 \cdot 10^0$	none	$< 2 \cdot 10^{-1}$
	CDHS	$3 \cdot 10^{-1}$	none	$< 5.3 \cdot 10^{-1}$

Only LSND demands non-zero  $U_{e4} U_{\mu4}$ ,  $U_{e5} U_{\mu5}$ , etc.

Is this consistent with upper limits derived from null short-baseline (NSBL) experiments?

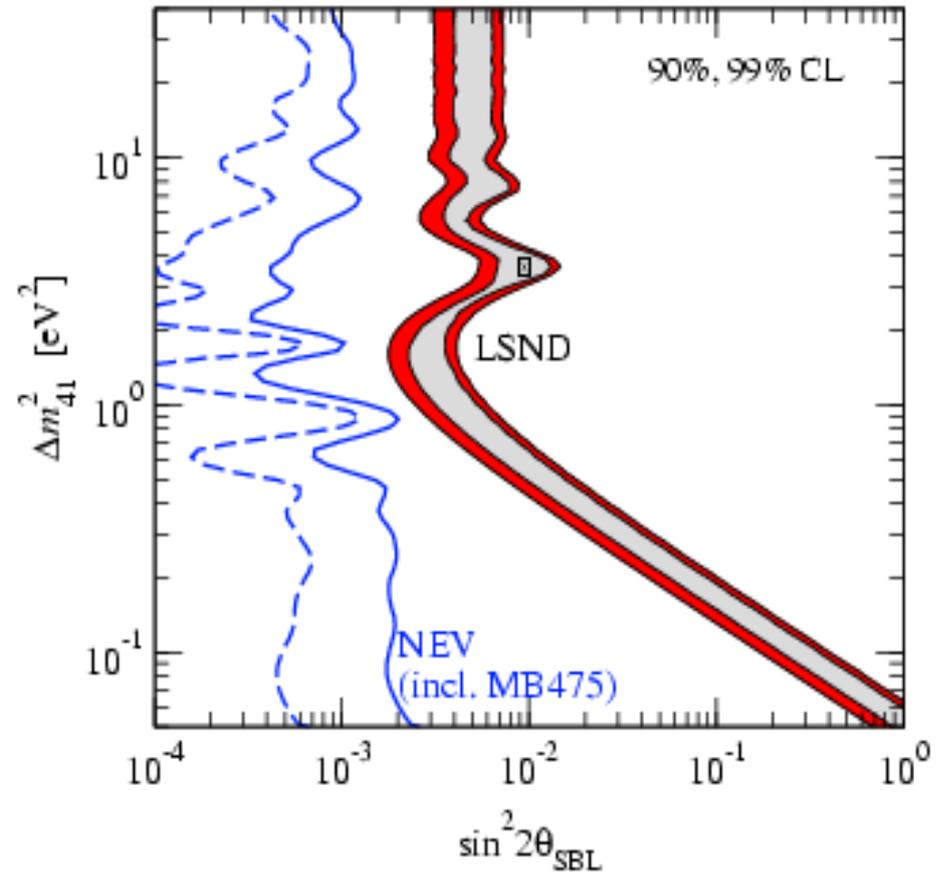
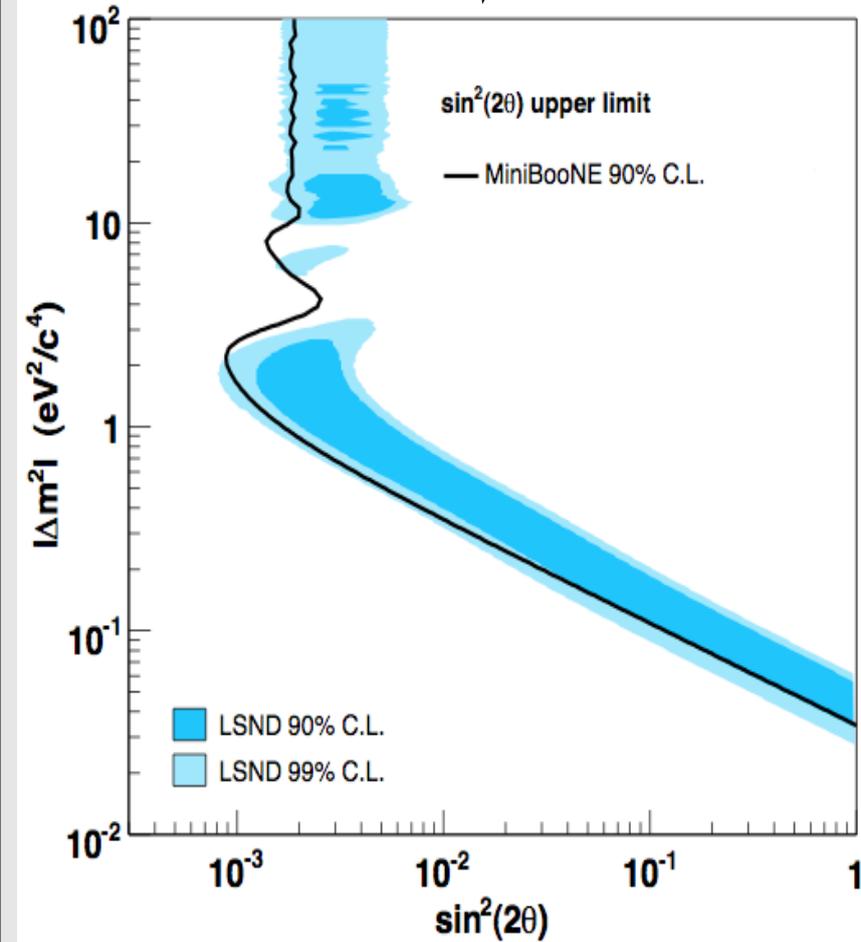
# First Results from MiniBooNE



MiniBooNE result **excludes** the LSND 90%CL allowed region **at > 90% CL...**

Counting Expt :  
 $300 \text{ MeV} < E^{QE} < 475 \text{ MeV}$   
**Excess over background : 3.7**

# MiniBooNE ( $E > 475 \text{ MeV}$ ) only, and (3+1) Global Fits



LSND ruled out with new MB results and within 3+1 models

# First Results from MiniBooNE

...MiniBooNE result assumes:

**CP-conserving, 2-neutrino oscillation scenario**

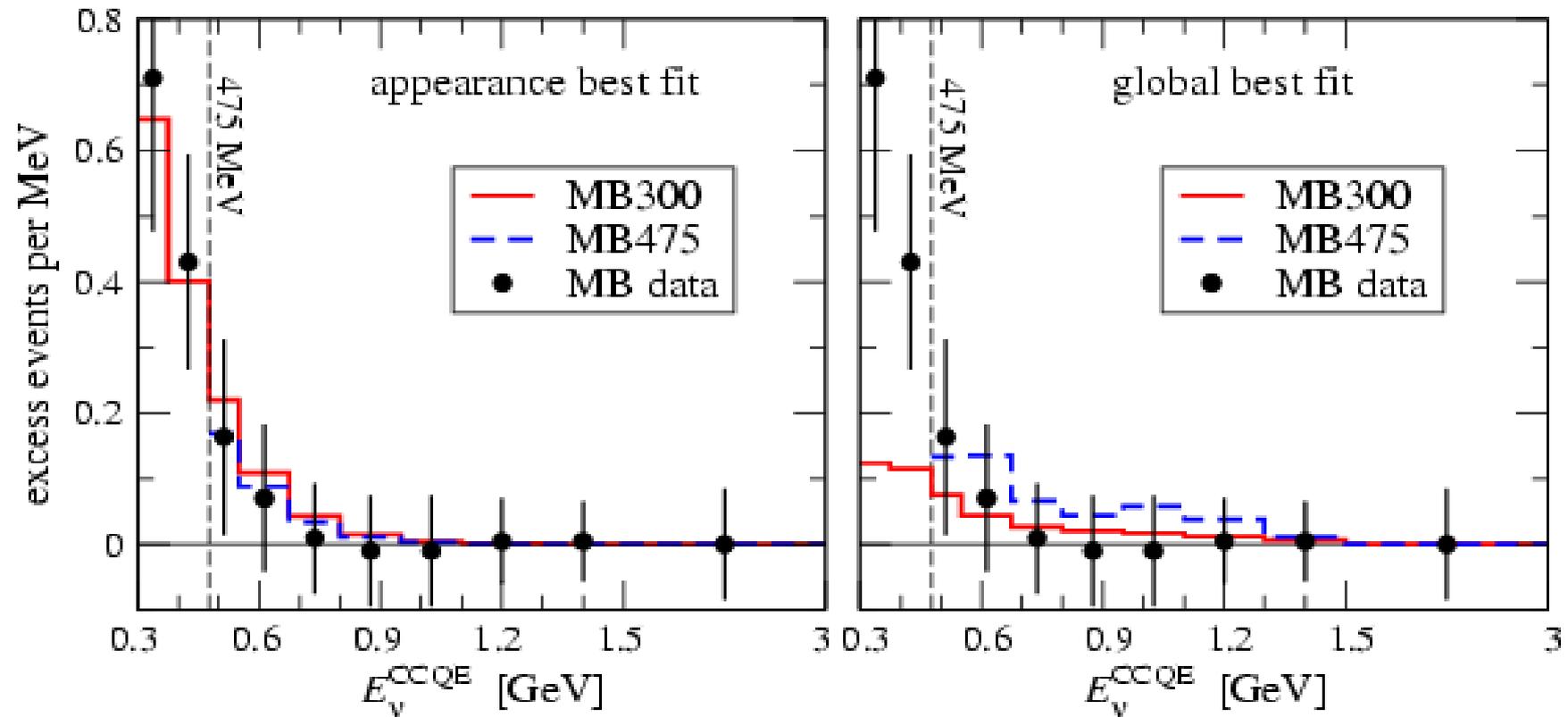
**$E > 475 \text{ MeV}$**

**Excess of  $\nu_e$  events at low energies:**

Currently investigating if this is a detector effect, or SM background...

Could be a manifestation of beyond the SM physics...

# 3+2 Global Fit Results



## 3+2 neutrino models:

- provide a **good fit** to LSND and the recent MB data
- can **account for the low energy event excess in MB**

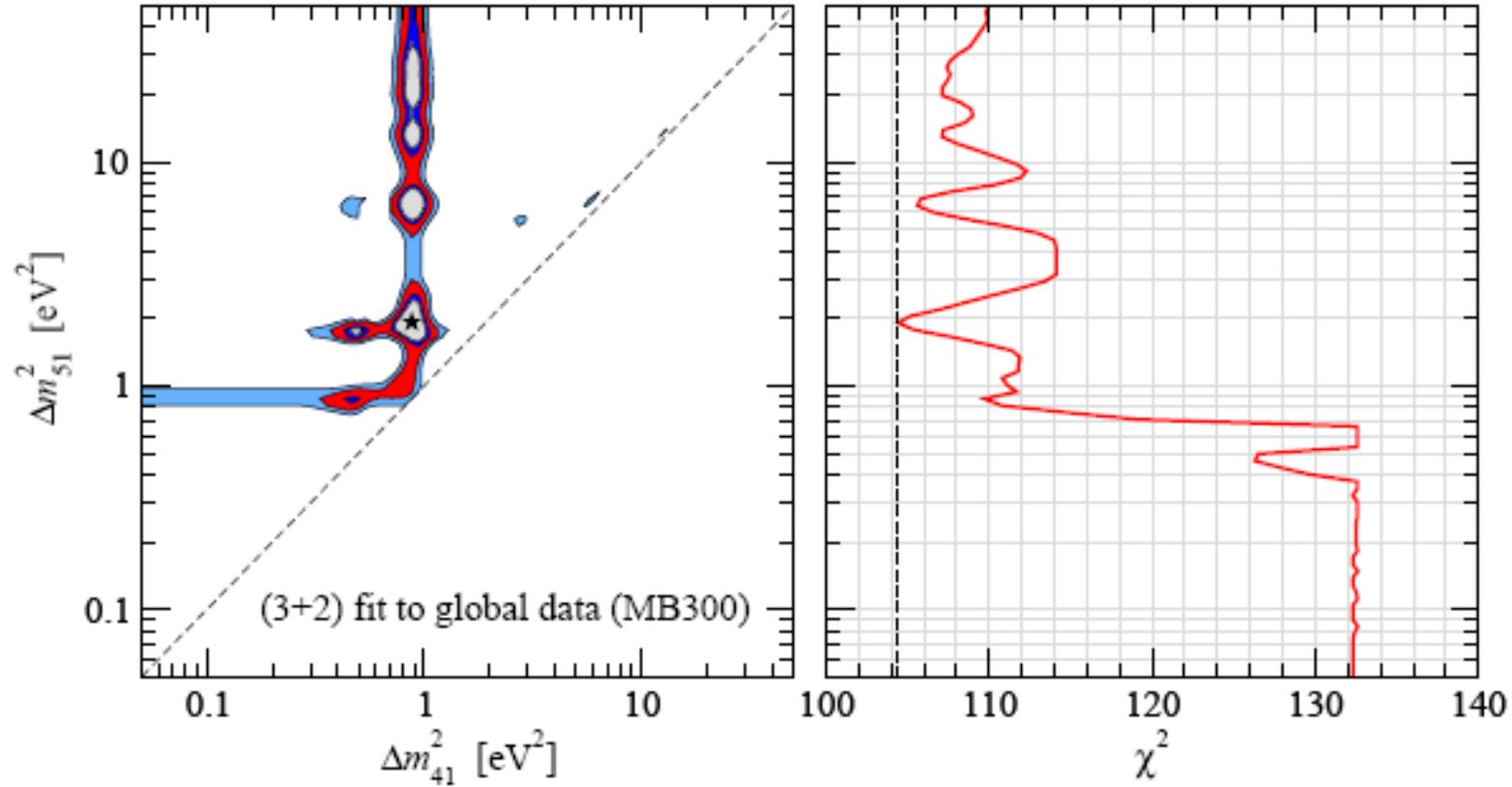
## However:

- there is significant **tension between appearance and disappearance data**

*Note: analysis done without full MiniBooNE error matrix*

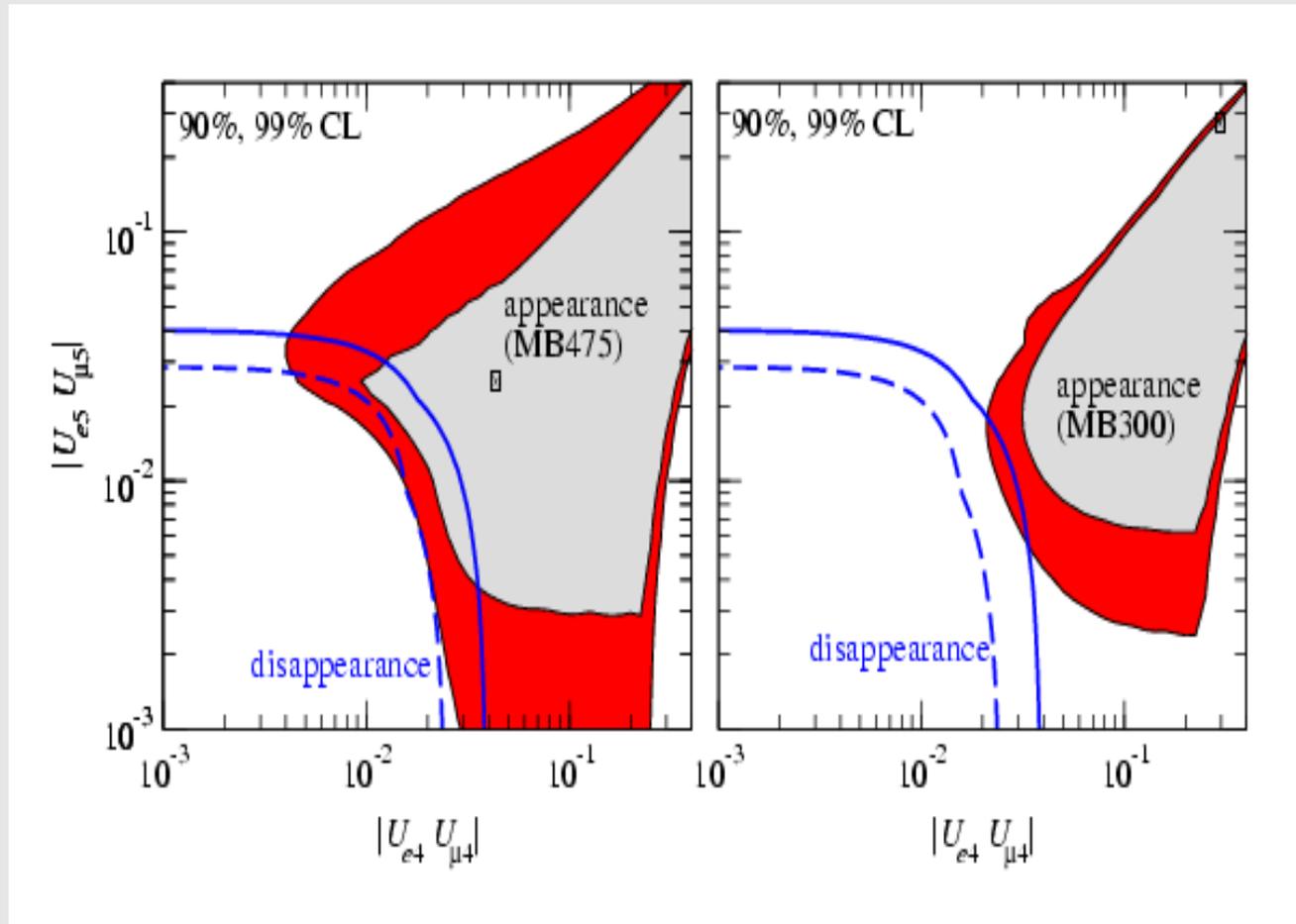
*MB will perform full analysis, G. Karagiorgi.*

## 3+2 Global Fit Results



data set	$ U_{e4}U_{\mu4} $	$\Delta m_{41}^2$	$ U_{e5}U_{\mu5} $	$\Delta m_{51}^2$	$\delta$	$\chi_{\min}^2/\text{dof}$	gof		
appearance (MB475)	0.044	0.66	0.022	1.44	$1.12\pi$	16.9/(29 - 5)	85%		
appearance (MB300)	0.31	0.66	0.27	0.76	$1.01\pi$	18.5/(31 - 5)	85%		
	$ U_{e4} $	$ U_{\mu4} $	$ U_{e5} $	$ U_{\mu5} $					
global data (MB475)	0.11	0.16	0.89	0.12	0.12	6.49	$1.64\pi$	94.5/(107 - 7)	63%
global data (MB300)	0.12	0.18	0.87	0.11	0.089	1.91	$1.44\pi$	104.4/(109 - 7)	41%

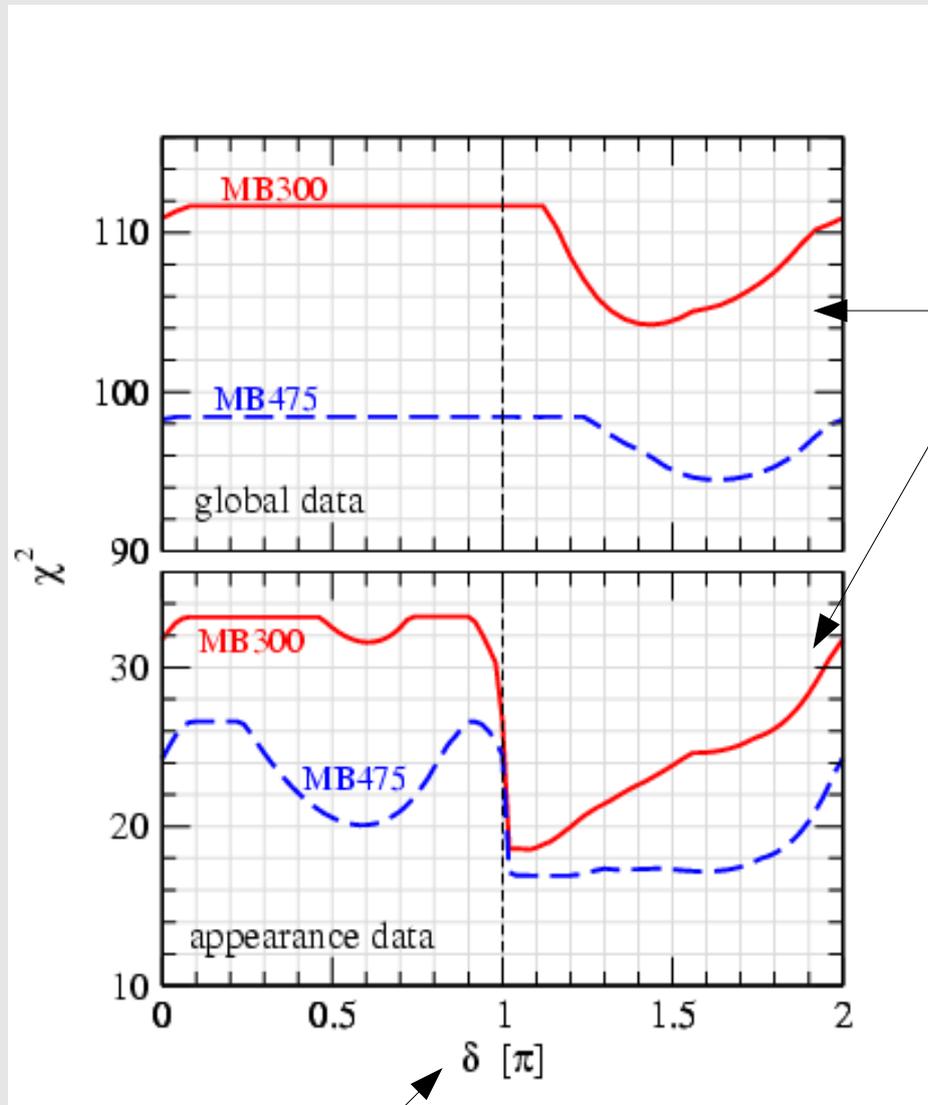
# 3+2 Tension with Disappearance Data



Large values of 3+2 appearance mixing amplitudes are inconsistent with disappearance amplitudes

# 3+2 Model CP Phase Fits

## Has CP violation been observed?



Improved fits with CP phase included!

MiniBooNE antineutrino oscillation analysis will be interesting.  
-could see LSND oscillations  
-may need more antineutrino data for decent sensitivity!

CP conserving value  $\delta = \pi$

# Conclusions

- Standard Model does not tell us a whole lot about sterile neutrino properties, only gives some general guidance,
  - Given neutrinos have mass, there are right handed singlets, but number of flavors and mass spectrum unknown.
  - They can have Dirac and/or Majorana mass terms.
  - See-saw mechanism may or may not be relevant.
- Sterile neutrinos have interesting properties/effects,
  - Sterile neutrinos may interact with extra D's.
  - Sterile neutrinos have measurable consequences in cosmology and astrophysics.
  - 3+2 models, with CP, can be constructed to fit SBL data.
- Up to future experiments to nail down neutrino properties,
  - Wouldn't it be nice to have a smoking gun measurement of active-sterile oscillations...